

> APPENDIX 1.0 THREE TIMESCALE

> restart; #causes the maple kernel to clear its internal memory so that Maple acts (almost) as if just started.

> deqn := diff(x(t), t, t) + x(t) + epsilon*(x(t)^2 - 1) * diff(x(t), t);
the equation of the van der Pol oscillator

$$deqn := \frac{d^2}{dt^2} x(t) + x(t) + \epsilon (x(t)^2 - 1) \left(\frac{d}{dt} x(t) \right) \quad (1)$$

> dsolve(deqn, x(t));
we are not able to solve the van der Pol equation in simple functions in maplesoft

$$x(t) = _a \text{ \&where } \left[\left\{ \left(\frac{d}{d_a} _b(_a) \right) _b(_a) + _a + \epsilon _b(_a) _a^2 - \epsilon _b(_a) = 0 \right\}, \left\{ _a = x(t), _b(_a) = \frac{d}{dt} x(t) \right\}, \left\{ t = \int \frac{1}{_b(_a)} d_a + _CI, x(t) = _a \right\} \right] \quad (2)$$

> xansatz := t → X[0](t, epsilon*t, epsilon^2*t) + epsilon*X[1](t, epsilon*t, epsilon^2*t) + epsilon^2*X[2](t, epsilon*t, epsilon^2*t); # perturbation series expansion

$$xansatz := t \rightarrow X_0(t, \epsilon t, \epsilon^2 t) + \epsilon X_1(t, \epsilon t, \epsilon^2 t) + \epsilon^2 X_2(t, \epsilon t, \epsilon^2 t) \quad (3)$$

> deqnseries := series(subs(epsilon^2*t = sigma, subs(epsilon*t = tau, eval(subs(x = xansatz, deqn))))), epsilon, 3);
substitutes equation (3) into equation (1) and collect like terms in powers of epsilon

$$\begin{aligned} deqnseries := & D_{1,1}(X_0)(t, \tau, \sigma) + X_0(t, \tau, \sigma) + (X_1(t, \tau, \sigma) + 2 D_{1,2}(X_0)(t, \tau, \sigma) \\ & + D_{1,1}(X_1)(t, \tau, \sigma) + (X_0(t, \tau, \sigma)^2 - 1) D_1(X_0)(t, \tau, \sigma)) \epsilon + (2 D_{1,3}(X_0)(t, \tau, \sigma) \\ & + D_{1,1}(X_2)(t, \tau, \sigma) + 2 D_{1,2}(X_1)(t, \tau, \sigma) + X_2(t, \tau, \sigma) + D_{2,2}(X_0)(t, \tau, \sigma) \\ & + (X_0(t, \tau, \sigma)^2 - 1) (D_2(X_0)(t, \tau, \sigma) + D_1(X_1)(t, \tau, \sigma)) + 2 X_0(t, \tau, \sigma) X_1(t, \tau, \\ & \sigma) D_1(X_0)(t, \tau, \sigma)) \epsilon^2 + O(\epsilon^3) \end{aligned} \quad (4)$$

> deqn0 := coeff(deqnseries, epsilon, 0); # lists the coefficients of the zeroth power of epsilon

$$deqn0 := D_{1,1}(X_0)(t, \tau, \sigma) + X_0(t, \tau, \sigma) \quad (5)$$

> deqn1 := coeff(deqnseries, epsilon, 1); # lists the coefficients of the first power of epsilon

$$\begin{aligned} deqn1 := & X_1(t, \tau, \sigma) + 2 D_{1,2}(X_0)(t, \tau, \sigma) + D_{1,1}(X_1)(t, \tau, \sigma) + (X_0(t, \tau, \sigma)^2 \\ & - 1) D_1(X_0)(t, \tau, \sigma) \end{aligned} \quad (6)$$

> deqn2 := coeff(deqnseries, epsilon, 2); # lists the coefficients of the second power of epsilon

$$\begin{aligned} deqn2 := & 2 D_{1,3}(X_0)(t, \tau, \sigma) + D_{1,1}(X_2)(t, \tau, \sigma) + 2 D_{1,2}(X_1)(t, \tau, \sigma) + X_2(t, \tau, \sigma) \\ & + D_{2,2}(X_0)(t, \tau, \sigma) + (X_0(t, \tau, \sigma)^2 - 1) (D_2(X_0)(t, \tau, \sigma) + D_1(X_1)(t, \tau, \sigma)) \\ & + 2 X_0(t, \tau, \sigma) X_1(t, \tau, \sigma) D_1(X_0)(t, \tau, \sigma) \end{aligned} \quad (7)$$

> deqn0; # we call the zeroth power coefficient

$$D_{1,1}(X_0)(t, \tau, \sigma) + X_0(t, \tau, \sigma) \quad (8)$$

> $dsolve(deqn0, X[0](t, \tau, \sigma));$ # we solve equation (8)

$$X_0(t, \tau, \sigma) = _F1(\tau, \sigma) \sin(t) + _F2(\tau, \sigma) \cos(t) \quad (9)$$

> $X[0] := (t, \tau, \sigma) \rightarrow A(\tau, \sigma) * \cos(t - \Phi(\tau, \sigma));$
 # we assume that the first linearized X_0 solution is a periodic solution with a given amplitude and period.

$$X_0 := (t, \tau, \sigma) \rightarrow A(\tau, \sigma) \cos(t - \Phi(\tau, \sigma)) \quad (10)$$

> $deqn0;$ # we show that equation (10) solves equation (8)

$$0 \quad (11)$$

> $deqn1temp := collect(combine(deqn1, trig), [\sin, \cos]);$ # we substitute for X_0 in equation (6) and obtain

$$deqn1temp := \left(-2 D_1(A)(\tau, \sigma) - \frac{1}{4} A(\tau, \sigma)^3 + A(\tau, \sigma) \right) \sin(t - \Phi(\tau, \sigma)) + X_1(t, \tau, \sigma) \quad (12)$$

$$+ 2 A(\tau, \sigma) \cos(t - \Phi(\tau, \sigma)) D_1(\Phi)(\tau, \sigma) + D_{1,1}(X_1)(t, \tau, \sigma) - \frac{1}{4} A(\tau, \sigma)^3 \sin(3t - 3\Phi(\tau, \sigma))$$

> $secular_terms := \{coeff(deqn1temp, \sin(t - \Phi(\tau, \sigma))), coeff(deqn1temp, \cos(t - \Phi(\tau, \sigma)))\}$ # we collect secular terms in equation (12)

$$secular_terms := \left\{ 2 A(\tau, \sigma) D_1(\Phi)(\tau, \sigma), -2 D_1(A)(\tau, \sigma) - \frac{1}{4} A(\tau, \sigma)^3 + A(\tau, \sigma) \right\} \quad (13)$$

> $dsolve(secular_terms, \{\Phi(\tau, \sigma), A(\tau, \sigma)\});$
 # we equate secular terms to zero and solve for Φ and A

$$\{A(\tau, \sigma) = 0, \Phi(\tau, \sigma) = \Phi(\tau, \sigma)\}, \left\{ A(\tau, \sigma) = -\frac{2}{\sqrt{1 + 4 e^{-\tau} _F1(\sigma)}}, \Phi(\tau, \sigma) = _F2(\sigma) \right\} \quad (14)$$

> $X[0] := (t, \tau, \sigma) \rightarrow 2 \cdot \cos(t - \Phi(\sigma));$ # as t tends to infinity, $A(\tau, \sigma)$ tends to the value 2; the solution $\Phi(\tau, \sigma) = _F2(\sigma)$, means Φ is a function of σ only.

$$X_0 := (t, \tau, \sigma) \rightarrow 2 \cos(t - \Phi(\sigma)) \quad (15)$$

> $deqn1temp := map(simplify, collect(combine(deqn1, trig), [\sin, \cos]));$
 # equation (12) simplifies to:

$$deqn1temp := X_1(t, \tau, \sigma) + D_{1,1}(X_1)(t, \tau, \sigma) - 2 \sin(3t - 3\Phi(\sigma)) \quad (16)$$

> $dsolve(deqn1temp, X[1](t, \tau, \sigma));$
 # as equation (16) has no secular terms, we next solve for the linearized solution, X_1

$$X_1(t, \tau, \sigma) = \sin(t) _F2(\tau, \sigma) + \cos(t) _F1(\tau, \sigma) - \frac{1}{4} \sin(3t - 3\Phi(\sigma)) \quad (17)$$

> $X[1] := (t, \tau, \sigma) \rightarrow -\frac{1}{4} \sin(3t - 3\Phi(\sigma));$
 # the second linearized X_1 solution is a periodic solution with amplitude and period and written as :

$$X_1 := (t, \tau, \sigma) \rightarrow -\frac{1}{4} \sin(3t - 3\Phi(\sigma)) \quad (18)$$

> $collect(combine(deqn1, trig), [\sin, \cos]);$ # we show that the linearised solutions X_0 and X_1 solves equation (6)

$$0 \quad (19)$$

> *deqn2temp* := map(simplify, collect(combine(*deqn2*, trig), [sin, cos]));
 # we substitute the values of X_0 and X_1 in equation (7) and obtain :

$$\begin{aligned} \text{deqn2temp} := & \frac{1}{4} (-1 + 16 D(\Phi)(\sigma)) \cos(t - \Phi(\sigma)) - \frac{5}{4} \cos(5t - 5\Phi(\sigma)) \\ & + D_{1,1}(X_2)(t, \tau, \sigma) + X_2(t, \tau, \sigma) - \frac{3}{4} \cos(3t - 3\Phi(\sigma)) \end{aligned} \quad (20)$$

> *secular_terms* := {coeff(*deqn2temp*, sin(t-Phi(sigma))), coeff(*deqn2temp*, cos(t-Phi(sigma)))}; # we collect secular terms as we did before

$$\text{secular_terms} := \left\{ 0, -\frac{1}{4} + 4 D(\Phi)(\sigma) \right\} \quad (21)$$

> *secular_terms* := 4*diff(Phi(sigma), sigma) = 1/4;

$$\text{secular_terms} := 4 \left(\frac{d}{d\sigma} \Phi(\sigma) \right) = \frac{1}{4} \quad (22)$$

> dsolve(*secular_terms*); # we solve for Φ in equation (22)

$$\Phi(\sigma) = \frac{1}{16} \sigma + _C1 \quad (23)$$

> Phi := sigma → $\frac{1}{16} \sigma + \text{phi}$; # we write Φ as a function of σ and ϕ

$$\Phi := \sigma \rightarrow \frac{1}{16} \sigma + \phi \quad (24)$$

> *X*[0](t, tau, sigma);

$$2 \cos\left(-t + \frac{1}{16} \sigma + \phi\right) \quad (25)$$

> *X*[1](t, tau, sigma);

$$\frac{1}{4} \sin\left(-3t + \frac{3}{16} \sigma + 3\phi\right) \quad (26)$$

> *deqn2temp* := map(factor, map(simplify, collect(combine(*deqn2*, trig), [sin, cos])));
 # substituting for Φ in equation (20) above, we have:

$$\begin{aligned} \text{deqn2temp} := & D_{1,1}(X_2)(t, \tau, \sigma) + X_2(t, \tau, \sigma) - \frac{5}{4} \cos\left(-5t + \frac{5}{16} \sigma + 5\phi\right) - \frac{3}{4} \cos\left(-3t \right. \\ & \left. + \frac{3}{16} \sigma + 3\phi\right) \end{aligned} \quad (27)$$

> dsolve(*deqn2temp*, *X*[2](t, tau, sigma)); # solving equation (27) for the linearized solution X_2 ,
 we have :

$$\begin{aligned} X_2(t, \tau, \sigma) = & \sin(t) _F2(\tau, \sigma) + \cos(t) _F1(\tau, \sigma) - \frac{5}{96} \cos\left(-5t + \frac{5}{16} \sigma + 5\phi\right) \\ & - \frac{3}{32} \cos\left(-3t + \frac{3}{16} \sigma + 3\phi\right) \end{aligned} \quad (28)$$

> *X*[2] := (t, tau, sigma) → $-\frac{5}{96} \cos\left(-5t + \frac{5}{16} \sigma + 5\phi\right) - \frac{3}{32} \cos\left(-3t + \frac{3}{16} \sigma + 3\phi\right)$;
 # the third linearized X_2 solution is a periodic solution with amplitude and period
 and written as

$$X_2 := (t, \tau, \sigma) \rightarrow -\frac{5}{96} \cos\left(-5t + \frac{5}{16} \sigma + 5\phi\right) - \frac{3}{32} \cos\left(-3t + \frac{3}{16} \sigma + 3\phi\right) \quad (29)$$

```
> collect(combine(deqn2, trig), [sin, cos]); # we show that substituting for  $X_2$ 
    in equation (27) solves that equation
                                0
(30)
```

```
> xansatz(t); # hence the solution or xansatz, as a function of t only is:
2 cos(-t + 1/16 ε2 t + φ) + 1/4 ε sin(-3 t + 3/16 ε2 t + 3 φ) + ε2 (-5/96 cos(-5 t + 5/16 ε2 t
+ 5 φ) - 3/32 cos(-3 t + 3/16 ε2 t + 3 φ))
(31)
```

```
> foo := dsolve([subs(epsilon=1/2, deqn), x(0)=2.0, D(x)(0)=0], range=0..100, numeric);
# numerical integration of the van der Pol with initial conditions
foo := proc(x_rkf45) ... end proc
(32)
```

```
> # tmp:=fsolve(subs(epsilon=1/2, {xansatz(0)=2.0})); # we set
epsilon=0.5 and solve for phi in equation (31) with initial
conditions
```

```
> tmp := {φ=0.2}
                                tmp := {φ=0.2}
(33)
```

```
> xapprox:=subs(epsilon=1/2, subs(tmp, xansatz(t))); # substituting
epsilon =0.5 and phi = .1396823540, the xansatz is
xapprox := 2 cos(-63/64 t + 0.2) + 1/8 sin(-189/64 t + 0.6) - 5/384 cos(-315/64 t + 1.0)
- 3/128 cos(-189/64 t + 0.6)
(34)
```

```
> with(plots): # we call the plot command
> pnum:=odeplot(foo, t=0..50, refine=1, color=red): # plot numerical
integration of the van der Pol in red
> pexp:=plot(xapprox, t=0..50, color=green): # plot the linearized
solutions in green
> display(pexp, pnum); # display both graphs
```

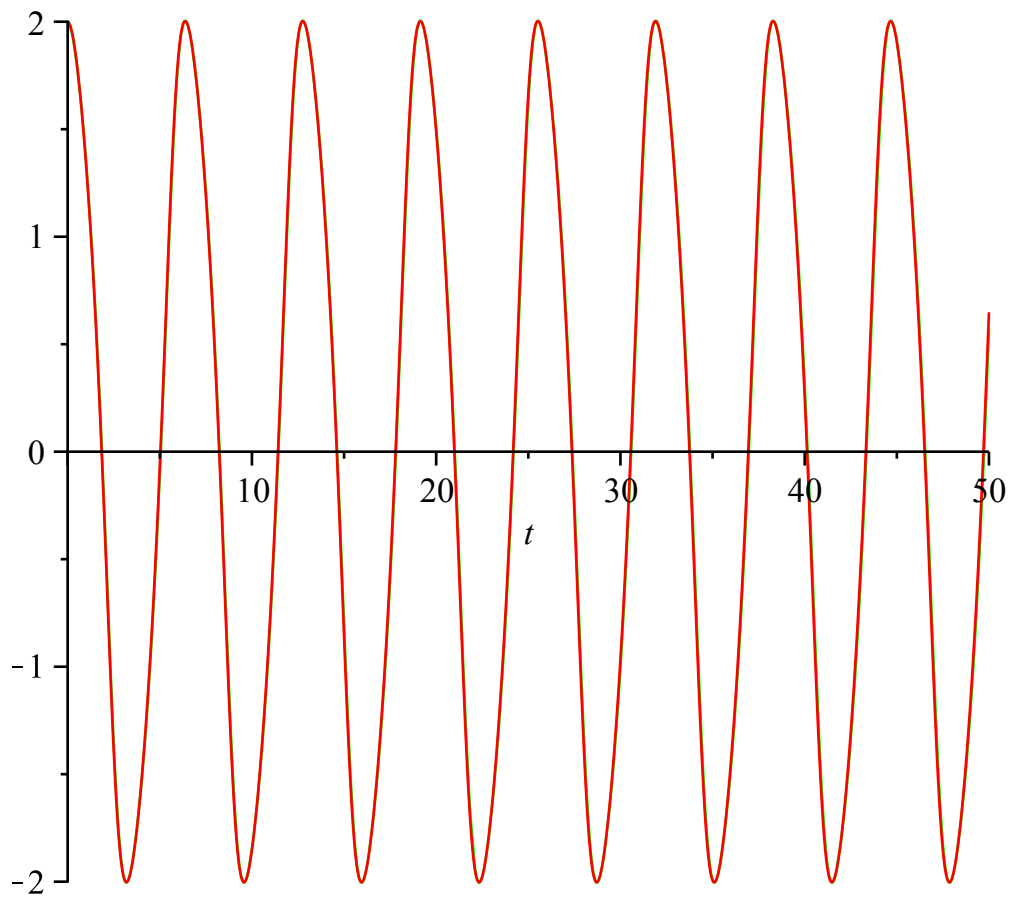


Figure 5.0(a) numerical integration of the van der Pol in red vs approximate solution in green by the 3-timing method for $\phi = 0.2$ and $\epsilon = 0.5$

```
> pnum2:=odeplot(foo,[x(t),D(x)(t)],t=0..10,refine=1,color=red):#
phase plot of numerical integration of the van der Pol in red
> pexp2:=plot([xapprox,diff(xapprox,t),t=0..10],color=green): #
phase plot of linearized solutions in green
> display(pnum2,pexp2); # display graphs
```

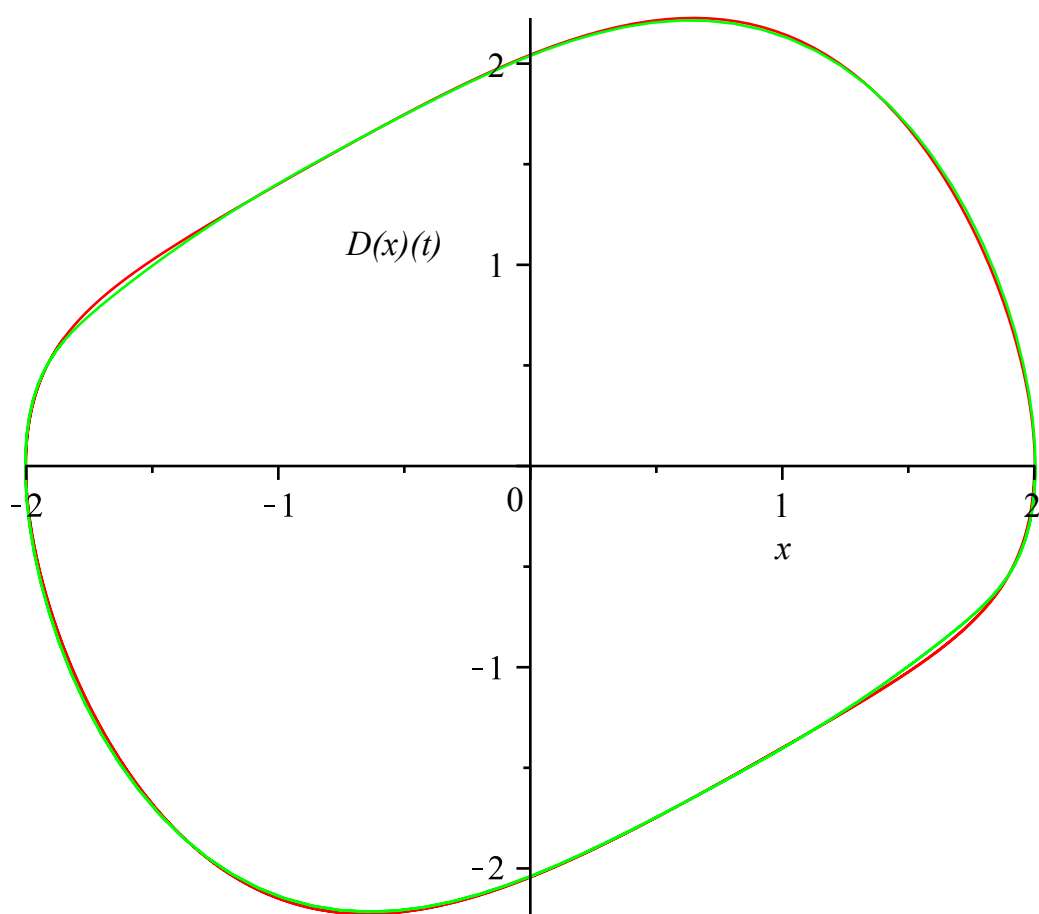


Figure 5.0(b) phase plot of the numerical integration of the van der Pol in red vs approximate solution in green by the 3-timing method for $\phi = 0.2$ and $\epsilon = 0.5$

